Is lattice simulation of two-dimensional super Yang-Mills possible?

KANAMORI Issaku



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Refs.

I.K. and H.Suzuki, Nucl.Phys.B in press [arXiv:0809.2856[hep-lat]] I.K., H.Suzuki and F.Suginio, Phys.Rev. D77 (2008), 091502 I.K., work in progress

Lattice Simulation of 2-dim super Yang-Mills is now possible

Introduction

Lattice Simulation of 2-dim super Yang-Mills is now possible

Plan

- 1. Introduction & Motivation
- 2. Formulation: Sugino model
- 3. What should we measure? Partially Conserved Super Current (PCSC)
- 4. Simulation result
- 5. Application: observing spontaneous SUSY breaking
- 6. Conclusion

















"Tell me why you guys choose physics in *English*"



the poor first victim

Lattice: non-perturbative regularization continuum limit ($\alpha \rightarrow 0$) in the end

successful in...

- QCD (gauge theory)
- gravity





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... treating fermion is not so easy what is fermions?

⇒putting supersymmetry = understanding fermions





Motivation

Lattice simulation: a non-perturbative method for field theory

- SUSY breaking Why our world is not supersymmetric?
- gauge/string duality tool for strong coupling gauge theory
- SUSY can be a symmetry of beyond Standard Model worth developing simulation techniques (will be found in LHC ?)
- "Experiment" for theoretical analysis poster by Suzuki

Formulation



Scalar Q for $N \ge 2$ on lattice





- topological twist \Rightarrow scalar Q on a site
- **simulation:** Catterall, Suzuki, Fukaya-I.K.-Suzuki-Takimi, I.K-Suzuki-Sugino



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Two types

1. Keep the whole SUSY (at finite *a*)

DKKN; Nagata; Arianos-D'Adda-Feo-Kawamoto-Saito; Kato-Sakamoto-So

2. partially at finite a, the whole is (automatically) restored in $a \rightarrow 0$ CKKU;Sugino



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We confirm this scenario in the simulation

Sugino, JHEP 01(2004)067

target: 2-dim $\mathcal{N} = (2, 2)$ SYM

nilpotent Q (Twisted) SUSY Algebra, continuum

$$Q^2 = \delta_{\phi}^{(\text{gauge})}$$
 $Q_0^2 = \delta_{\overline{\phi}}^{(\text{gauge})}$ $\{Q, Q_0\} = 2i\partial_0 + 2\delta_{A_0}^{(\text{gauge})}$

Action (dimensional reduction from 4-dim $\mathcal{N} = 1$)

$$S = \frac{1}{g^2} \int d^2 x \operatorname{tr} \left\{ \frac{1}{2} F_{MN} F_{MN} + \Psi^T C \Gamma_M D_M \Psi + \hat{H}^2 \right\}$$

$$\Psi^T = (\psi_0, \psi_1, \chi, \eta/2) \quad (\text{under a suitable rep. of } \Gamma_M)$$

$$A_M = (\text{gauge field, scalar})$$

$$\hat{H} = \text{aux. field}$$

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Q-exact action (continuum)

$$S = Q(...) = \frac{1}{g^2} \int d^2 x \, \text{tr} \Big\{ \frac{1}{4} F_{01}^2 + D_\mu \phi D_\mu \overline{\phi} + \frac{1}{4} [\phi, \overline{\phi}]^2 + i \psi_\mu D_\mu \eta + \dots - \frac{1}{4} \eta [\phi, \eta] + \dots \Big\}$$

$$QA_{\mu} = \psi_{\mu} \qquad Q\psi_{\mu} = iD_{\mu}\phi \qquad Q_{0}A_{0} = \frac{i}{2}\eta \quad Q_{0}\eta = -2iD_{0}\overline{\phi}$$
$$Q\phi = 0 \qquad \dots \qquad Q_{0}A_{1} = -\chi \quad Q_{0}\chi = iD_{1}\overline{\phi}$$
$$Q_{0}\overline{\phi} = 0 \qquad \dots$$

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target: 2-dim $\mathcal{N} = (2, 2)$ SYM nilpotent Q Lattice version

$$Q^2 = \delta_{\phi}^{(gauge)}$$

Q-exact action (lattice)

$$S = Q(\dots) = S[U(x, \mu), \phi(x), \overline{\phi}(x), H(x)$$
bosons

$$\eta(x), \chi(x), \psi_0(x), \psi_1(x)]$$
fermions

$$QU(x, \mu) = i\psi_{\mu}(x)U(x, \mu)$$

$$Q\psi_{\mu}(x) = i\psi_{\mu}(x)\psi_{\mu}(x)$$

$$-i(\phi(x) - U(x, \mu)\phi(x + \hat{\mu})U(x, \mu)^{-1})$$

$$Q\phi = 0$$

$$U$$

$$U$$

Symmetries

(continuum) \Rightarrow (lattice)

- 4 supercharges: $\{Q_0, Q_1, \tilde{Q}, Q\} \Rightarrow Q$
- R-symmetry: $U(1)_A, U(1)_V, Z_2 \Rightarrow U(1)_A$

($U(1)_A$: a rotational symmetry on 2-3 plane)

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- Other part of SUSY: broken at $a \neq 0$, but will be restored in $a \rightarrow 0$

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We will confirm this scenario

What should we measure? to check the restoration of SUSY

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Possible problem

• perturbative power counting: non-pertubativitly...?

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Powerful non-perturbative tool

• Simulation! : conservation of "super current"

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- 2. scalar mass term
- 3. boundary condition: anti-periodic in temporal direction for fermion (thermal)

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PCSC relation					
(Separate	the	effect	of	lattice	artifact)
satisfied	\Rightarrow	the latt	ice a	artifact v	vanishes
not satisfied	does not vanish				

PCSC relation

4 supercharges: $Q_A = \{Q_0, Q_1, \tilde{Q}, Q\}$

Partially conserved supercurrent:

$$\partial_{\mu} \mathcal{J}^{A}_{\mu} = 0 \Rightarrow \partial_{\mu} \mathcal{J}^{A}_{\mu} = \mu^{2}/g^{2} Y^{A}$$
 (PCSC) μ : scalar mass

$$\langle \partial_{\mu} \mathcal{J}^{A}_{\mu}(x) X^{A}(0) \rangle - \frac{\mu^{2}}{g^{2}} \langle Y^{A}(x) X^{A}(0) \rangle = -i\delta^{2}(x) \langle Q^{A} X^{A}(0) \rangle$$

measure
$$\frac{\langle \partial_{\mu} \mathcal{J}^{A}_{\mu}(x) X^{A}(0) \rangle}{\langle Y^{A}(x) X^{A}(0) \rangle}$$
 as a function of μ^{2}/g^{2}

$$Y^{A} = -2[C(\Gamma_{2} \operatorname{tr}(A_{2}\Psi) + \Gamma_{3} \operatorname{tr}(A_{3}\Psi))]^{A}$$

~ (scalar) × (fermion)

 $X^{A} = \frac{1}{g^{2}} [\Gamma_{0}(\Gamma_{2} \operatorname{tr}(A_{2} \Psi) + \Gamma_{3} \operatorname{tr}(A_{3} \Psi))]^{A}$

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Furthermore,

- no operator mixing for \mathcal{J} in 1-loop
- higher loops are suppressed because of super renormalizability

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Simple analysis is OK

Simulation Result

PCSC relation

$$\langle \partial_{\mu} \mathcal{J}^{A}_{\mu}(x) X^{A}(0) \rangle - \frac{\mu^{2}}{g^{2}} \langle Y^{A}(x) X^{A}(0) \rangle = -i\delta^{2}(x) \langle Q^{A} X^{A}(0) \rangle$$



 20×10 , ag = 0.1414, input $\mu^2/g^2 = 1.0$, \mathcal{J} for Q_0

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PCSC relation (continuum limit)



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PCSC is satisfied \Rightarrow no SUSY breaking due to lattice artifact

massless case: never thermalized

because of the flat direction in the potential



anti-PBC \Rightarrow PBC: Noisy



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Simulation detail

- Computer: Riken Super Combined Cluster (RSCC)
- Algorithm: Rational Hybrid Monte Carlo (RHMC)

+ Multi-time step acceleration

- lattice size: 12 × 6–30 × 10
- *ag* = 0.2357, 0.2, 0.1768, 0.1414
- 800–1,800 configurations

Applications

PCSC relation \Rightarrow lattice simulation works

Application

• Observing spontaneous SUSY breaking

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... before that, flat direction?

(How to "define" the theory)

- Define as an extrapolation of $\mu \rightarrow 0$
- Small mass: $\mu < 1/L$
- large N_C
- (unstable)

PCSC relation \Rightarrow lattice simulation works

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flat direction? (How to "define" the theory)

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cf. I.K.-Suzuki-Sugino

 $(\{Q, Q_0\} = 2i\partial_0)$

- order parameter $\langle \mathcal{H} \rangle = \frac{1}{2} \langle Q \mathcal{J}_0^{(0)} \rangle$
- measure under thermal boundary condition
- extrapolate the scalar mass to zero
- extrapolate to zero temperature: ground state energy density *E*

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 at zero temperature $\begin{cases} = 0 & \text{SUSY} \\ \neq 0 & \text{SUSY} \end{cases}$

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Cf. No analytic result for this system (maybe broken? a conjecture by Hori-Tong)

not broken?

Ground state energy is consistent with 0

Conclusion

Lattice simulation of 2-dim super Yang-Mills is now possible Formulation with 1 nilpotent Q \Rightarrow whole SUSY is restored in the cont. limit (no lat. artifact)

- simulation: 2-dim $\mathcal{N} = (2, 2)$ SYM, SU(2) (Sugino model)
- scalar mass term ⇒ Partially Conserved Super Current (PCSC) relation

Application

- Observing (no) spontaneous SUSY breaking
- : (Poster by Suzuki)

Thank you.